

Bianchi Type V Inflationary Universe with Decaying Vacuum Energy (Λ)

¹Raj Bali and ²Swati Singh

¹Department of Mathematics, CSIR Emeritus Scientist,
University of Rajasthan, Jaipur - 302 004 India

²Department of Mathematics, University of Rajasthan, Jaipur - 302 004 India

Abstract: Bianchi Type V cosmological inflationary model with flat potential and decaying vacuum energy is investigated. To get the deterministic solution in terms of cosmic time t , we assume that the decaying vacuum energy $\Lambda \sim \frac{\alpha}{R^2}$ where R is scale factor as considered by Chen and Wu [38], α being a constant, we find that the spatial volume increases exponentially indicating the inflationary scenario in the model. The model represents decelerating and accelerating phases both which matches with the recent astronomical observations. The anisotropy is maintained throughout in the model. However, if the constant $L = 0$ then the model isotropizes. The rate of Higg's field decreases slowly with time. The model has Point Type Singularity at $T = 0$ [39].

Key words: Bianchi V • Inflationary universe • Vacuum energy (Λ)

INTRODUCTION

Bianchi models are significant in the study because these models are homogeneous and anisotropic from which the process of isotropization of the universe is studied through the passage of time. The study of Bianchi Type V cosmological models create more interest in the study because these models are anisotropic generalization of open FRW models and allow arbitrarily small anisotropy levels at any constant of cosmic time. Bianchi Type V models have been studied in detail by number of authors viz. Farnsworth [1], Roy and Singh [2], Banerjee and Sanyal [3], Maartens and Nel [4], Collins [5], Wainwright *et al.* [6], Coley [7], Bali and Meena [8], Bali and Kumawat [9].

Inflationary universes create more interest in the study because these universes play a significant role in solving number of outstanding problems in cosmology like homogeneity, the isotropy, the horizons, flatness and primordial monopole problem in grand unified field theories. Guth [10] introduced the concept of inflation while investing the problem of why we do not see magnetic monopole today. He found that a positive-energy false vacuum generates an exponential expansion of space in general relativity. In particular, our universe is

homogeneous and isotropic to a very high degree of precision. Such a universe is described by FRW space-time. Several version of inflationary scenario are studied by number of authors viz. Linde [11], Wald [12], Barrow [13], Burd and Barrow [14], La and Steinhardt [15] in FRW space-time. Rothman and Ellis [16] have pointed out that we can have a solution of isotropic problem if we work with anisotropic metric and these metrics can be isotropized under very general circumstances. Stein-Schabes [17] has shown that inflation will take place if effective potential $V(\phi)$ has flat region while Higg's field evolves slowly but the universe expands in an exponential way due to vacuum field energy. Therefore, it is interesting to investigate inflationary scenario in anisotropic metric which isotropizes at late time or in a very general circumstances. Keeping such type of investigations, Bali and Jain [18] investigated inflationary scenario in LRS Bianchi Type I space-time in the presence of massless scalar field with flat potential. Recently Bali [19] investigated inflationary scenario in Bianchi Type I space-time with flat potential considering the scale factor $a \sim e^{3Ht}$. In this paper, we have discussed inflationary scenario in Bianchi Type V space-time with flat potential and decaying vacuum energy (Λ).

The cosmological constant (Λ) was introduced by Einstein to find the solution of static universes because at that time universe was supposed to be static. But after the discovery of Hubble constant, it was realized that universe is expanding. Also FRW obtained an expanding dust filled homogeneous and isotropic model in which there was no need to introduce the cosmological constant (Λ) into the Einstein's field equations. Einstein rejected the introduction of Λ term into his field equations after the realization that universe is expanding. However, two independent groups led by Riess *et al.* [20] and Perlmutter *et al.* [21] used Type Ia Supernovae and showed that universe is accelerating. This discovery provided the first direct evidence that Λ is non-zero with $\Lambda \sim 1.7 \times 10^{8121}$ Planck units. It is now commonly believed by Scientific community that via the cosmological constant, a kind of repulsive pressure dubbed as dark energy, is the most suitable candidate to explain recent observations that universe appears to be expanding and accelerating. According to the first year data not of Supernovae Legacy Survey (SNLS), dark energy behaves like the cosmological constant to a precision of 10% [22]. Λ CDM models agree closely with almost all the established cosmological abbreviations. Obviously, this extremely small value of cosmological constant, indicates that vacuum energy density (Λ) is not a strict constant but decays as the universe expands. A wide range of observations suggest that the cosmological constant Λ is the most favourable candidate of dark energy representing energy density of vacuum. Recently Barrow and Shaw [23] suggested that cosmological constant term corresponds to a very small value of the order of 10^{-122} when applied to Friedmann universe. A number of cosmological models in which Λ decays with time have been investigated by several authors viz. Bertolami [24], Ram [25], Berman [26], Beesham [27], Abdussattar and Vishwakarma [28], Sahni and Starobinski [29], Bronnikov *et al.* [30], Singh and Chaubey [31], Singh *et al.* [32], Bali and Singh [33,34], Ram and Verma [35] Abdussattar and Prajapati [36], Bali *et al.* [37].

In this paper, we have investigated inflationary scenario in Bianchi Type V space-time with flat potential and decaying vacuum energy density (Λ). We find that spatial volume increases exponentially indicating inflationary scenario in the model. The vacuum energy density Λ decreases with time. The model describes a unified expansion history of the universe indicating decelerating and accelerating phases both.

The anisotropy is maintained throughout. However, if the constant $L = 0$ then the model isotropizes. The Higg's field evolves slowly but universe expands.

Metric and Field Equations: We consider Bianchi Type V line-element in orthogonal form as

$$ds^2 = -dt^2 + A^2(t) dx^2 + e^{2x} (B^2(t)dy^2 + C^2(t) dz^2) \quad (1)$$

where A, B, C are metric potentials.

We assume the coordinates to be comoving so that $v^1 = 0 = v^2 = v^3, v^4 = 1$. The Einstein's field equations (in gravitational units $8\pi G = c = 1$) with time varying cosmological term $\Lambda(t)$ are given by

$$R_{ij} - \frac{1}{2}R g_{ij} + \Lambda g_{ij} = -T_{ij} \quad (2)$$

with

$$T_{ij} = \partial_i \phi \partial_j \phi - \left[\frac{1}{2} \partial_\rho \phi \partial^\rho \phi + V(\phi) \right] g_{ij} \quad (3)$$

and

$$\frac{1}{\sqrt{-g}} \partial_\mu (\sqrt{-g} \partial^\mu \phi) = -\frac{dV}{d\phi} \quad (4)$$

The field equations (2) for the line-element (1) lead to

$$\frac{B_{44}}{B} + \frac{C_{44}}{C} + \frac{B_4 C_4}{BC} - \frac{1}{A^2} + \Lambda(t) = -\frac{1}{2} \dot{\phi}^2 + V(\phi) \quad (5)$$

$$\frac{A_{44}}{A} + \frac{C_{44}}{C} + \frac{A_4 C_4}{AC} - \frac{1}{A^2} + \Lambda(t) = -\frac{1}{2} \dot{\phi}^2 + V(\phi) \quad (6)$$

$$\frac{A_{44}}{A} + \frac{B_{44}}{B} + \frac{A_4 B_4}{AB} - \frac{1}{A^2} + \Lambda(t) = -\frac{1}{2} \dot{\phi}^2 + V(\phi) \quad (7)$$

$$\frac{A_4 B_4}{AB} + \frac{A_4 C_4}{AC} + \frac{B_4 C_4}{BC} - \frac{3}{A^2} + \Lambda(t) = \frac{1}{2} \dot{\phi}^2 + V(\phi) \quad (8)$$

$$\frac{2A_4}{A} - \frac{B_4}{B} - \frac{C_4}{C} = 0 \quad (9)$$

The equation for scalar field (4) leads to

$$\ddot{\phi} + \left(\frac{A_4}{A} + \frac{B_4}{B} + \frac{C_4}{C} \right) \dot{\phi} = \frac{dV}{d\phi} \quad (10)$$

Solution of Field Equations: We are interested in inflationary solution so flat region is considered. Thus $V(\phi)$ is constant. Now equation (10) leads to

$$\ddot{\phi} + \left(\frac{A_4}{A} + \frac{B_4}{B} + \frac{C_4}{C} \right) \dot{\phi} = 0 \quad (11)$$

From equation (11), we have

$$\dot{\phi} = \frac{l}{ABC} \quad (12)$$

where l is constant of integration.

The scale factor R is given by

$$R^3 = ABC = A^3 \quad (13)$$

as $BC = A^2$ from equation (9). Equations (5) and (8) lead to

$$\frac{B_{44}}{B} + \frac{C_{44}}{C} + \frac{A_4 B_4}{AB} + \frac{A_4 C_4}{AC} + \frac{B_4 C_4}{BC} - \frac{4}{A^2} + 2\Lambda = 2k \quad (14)$$

where $V(\phi) = \text{constant} = k$. To get the deterministic solution, we assume that $\Lambda = \frac{\alpha}{R^2}$ as considered by Chen and Wu (1990).

For the sake of simplicity, we take

$$\alpha = 2 \quad (15)$$

Thus equations (14) and (15) lead to

$$\frac{B_{44}}{B} + \frac{C_{44}}{C} + \frac{2B_4 C_4}{BC} + \frac{A_4 B_4}{AB} + \frac{A_4 C_4}{AC} = 2k \quad (16)$$

Using equation (9) in equation (16), we have

$$\frac{B_{44}}{B} + \frac{C_{44}}{C} + \frac{1}{2} \left(\frac{B_4}{B} + \frac{C_4}{C} \right)^2 + \frac{2B_4 C_4}{BC} = 2k \quad (17)$$

Equations (6) and (7) lead to

$$\frac{B_{44}}{B} - \frac{C_{44}}{C} + \frac{A_4}{A} \left(\frac{B_4}{B} - \frac{C_4}{C} \right) = 0 \quad (18)$$

which again leads to

$$\frac{(CB_4 - BC_4)_4}{(CB_4 - BC_4)} = -\frac{A_4}{A} = -\frac{1}{2} \left(\frac{B_4}{B} + \frac{C_4}{C} \right) \quad (19)$$

using equation (9)

From equation (19), we have

$$C^2 \left(\frac{B}{C} \right)_4 = \frac{L}{\mu^{1/2}} \quad (20)$$

where L is constant of integration. To find the solution of equation (17) and (20), we assume that $BC = \mu$ and $\frac{B}{C} = v$.

Thus equations (17) and (20) lead to

$$2\mu_{44} + \frac{1}{\mu} \mu_4^2 = 4k \mu \quad (21)$$

and

$$\frac{v_4}{v} = \frac{L}{\mu^{3/2}} \quad (22)$$

Equation (21) leads to

$$f^2 = \frac{4}{3} k \mu^2 + \frac{\gamma}{\mu} \quad (23)$$

where $\mu_4 = f(\mu), \mu_{44} = f f^1, f^1 = df/d\mu$.

Equation (23) leads to

$$\frac{\sqrt{\mu} du}{\sqrt{(\mu^{3/2})^2 + \gamma/\beta^2}} = \beta dt \quad (24)$$

where $\frac{4k}{3} = \beta^2$ and $\gamma_1^2 = \frac{\gamma}{\beta^2}$. From equation (24), we have

$$\mu^{3/2} = \gamma_1 \sinh \left(\frac{3\beta}{2} t + \gamma_2 \right) \quad (25)$$

γ_2 being constant of integration. Now equations (22) and (25) lead to

$$v = \gamma_3 \tan h^{\frac{2L}{3\beta\gamma_1}} \left(\frac{T}{2} \right) \quad (26)$$

where $\frac{3\beta}{2}t + \gamma_2 = T$ and γ_3 being constant of integration.

Thus

$$A^2 = BC = \mu = \gamma_1^{2/3} \sinh^{2/3} T \quad (27)$$

$$B^2 = \mu v = \gamma_3 (2\gamma_1)^{2/3} \sinh^{\left(\frac{2}{3} + \frac{2L}{3\beta\gamma_1}\right)} \left(\frac{T}{2} \right) \cosh^{\left(\frac{2}{3} - \frac{2L}{3\beta\gamma_1}\right)} \left(\frac{T}{2} \right) \quad (28)$$

$$C^2 = \frac{\mu}{v} = \frac{(2\gamma_1)^{2/3}}{\gamma_3} \sinh^{\left(\frac{2}{3} - \frac{2L}{3\beta\gamma_1}\right)} \left(\frac{T}{2} \right) \cosh^{\left(\frac{2}{3} + \frac{2L}{3\beta\gamma_1}\right)} \left(\frac{T}{2} \right) \quad (29)$$

After suitable transformation of coordinates, the metric (1) leads to the form

$$ds^2 = -\frac{4}{9\beta^2} dT^2 + \sinh^{2/3} T dX^2 + e^{\frac{2X}{\gamma_1^{1/3}}} \left[\sinh^{\left(\frac{2}{3} + \frac{2L}{3\beta\gamma_1}\right)} \frac{T}{2} \cosh^{\left(\frac{2}{3} - \frac{2L}{3\beta\gamma_1}\right)} \frac{T}{2} dY^2 + \sinh^{\left(\frac{2}{3} - \frac{2L}{3\beta\gamma_1}\right)} \frac{T}{2} \cosh^{\left(\frac{2}{3} + \frac{2L}{3\beta\gamma_1}\right)} \frac{T}{2} dZ^2 \right] \quad (30)$$

where

$$\gamma_1^{1/3} x = X, \quad \gamma_3^{1/2} (2\gamma_1)^{1/3} y = Y, \quad \frac{(2\gamma_1)^{1/3}}{\gamma_3^{1/2}} z = Z$$

Some Physical and Geometrical Features: The vacuum energy density (Λ) is given by

$$\Lambda = \frac{2}{R^2} = \frac{2}{A^2} = \frac{2}{\gamma_1^{2/3} \sinh^{2/3} T} \quad (31)$$

The rate of Higg's field (ϕ) is given by equation (12) as

$$\dot{\phi} = \frac{\ell}{A^3} = \frac{\ell}{\gamma_1 \sinh T} \quad (32)$$

which leads to

$$\phi = \frac{2\ell}{3\beta\gamma_1} \log \tanh \frac{T}{2} + N \quad (33)$$

where N is constant of integration.

The spatial volume (R^3) for the model (30) is given by

$$R^3 = \gamma_1 \sinh T \quad (34)$$

The expansion (θ), shear (σ), the deceleration parameter (q) are given by

$$\begin{aligned} \theta &= \frac{A_4}{A} + \frac{B_4}{B} + \frac{C_4}{C} = \frac{3}{2} \left(\frac{B_4}{B} + \frac{C_4}{C} \right) = \frac{3\mu_4}{2\mu} \\ &= \frac{3\beta}{2} \coth T = \frac{3\beta}{2} \left(\frac{e^T + e^{-T}}{e^T - e^{-T}} \right) \end{aligned} \quad (35)$$

$$\sigma = \frac{1}{2} \left(\frac{B_4}{B} - \frac{C_4}{C} \right) = \frac{1}{2} \frac{v_4}{v} = \frac{L}{2\gamma_1 \sinh T} = \frac{L}{\gamma_1 (e^T - e^{-T})} \quad (36)$$

$$q = -\frac{\ddot{R}/R}{\dot{R}^2/R^2} = -3 \tan^2 T + 2 \quad (37)$$

DISCUSSION

We observe that expansion (θ), shear (σ) and decaying vacuum energy density (Λ) all diverge at $T = 0$. The model (30) starts with a big-bang from its singular state at $T = 0$ and continues to expand till $T = \infty$. The spatial volume (R^3) increases exponentially as T increases. Thus the model represents inflationary scenario. For large values of T , $\frac{\sigma}{\theta} \rightarrow 0$ which implies that

the model approaches isotropy at late times. When $T \rightarrow 0$ then $\frac{\sigma}{\theta}$ is finite. Hence the model represents anisotropic

space-time initially and isotropizes at late times. The model describes a unified expansion history of the universe which starts with decelerating expansion and the expansion accelerates at late time. The decelerating expansion at initial epoch provides obvious provision for the formation of large structure in the universe. The formation of structure is better supported by decelerating expansion. Thus the model is astrophysically relevant.

Also late time acceleration is in agreement with the observations of 16 type Ia supernovae made by Hubble space Telescope (HST)(Riess *et al.* [40]). The decaying vacuum energy density (Λ) is initially large but decreases as time passes. This result matches with astronomical observations.

CONCLUSION

The spatial volume (R^3) increases exponentially with time representing inflationary scenario. The model represents anisotropic space time initially but isotropizes at late time. The model describes unified expansion history of universe which starts with decelerating expansion and the accelerating expansion at late time. Also at late time, acceleration is in agreement with the observation of 16 Type Ia Supernovae made by Hubble space telescope. The vacuum energy (Λ) is initially large but decreases with time. This result matches with astronomical observations. Thus inflationary scenario exists with decaying vacuum energy in Bianchi Type V space-time.

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